

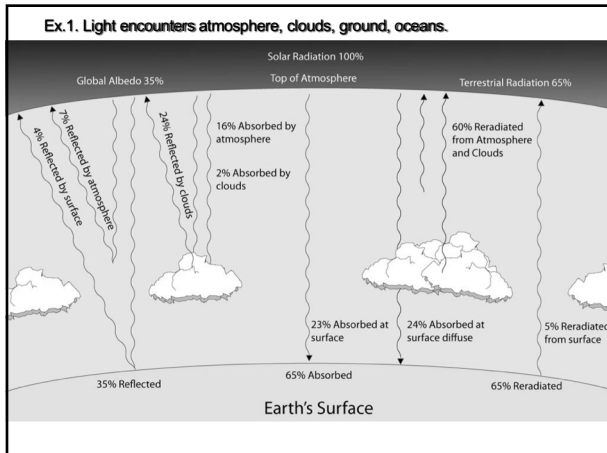


Wave Incidence

[Chapter 10 cont, Saḍiku]

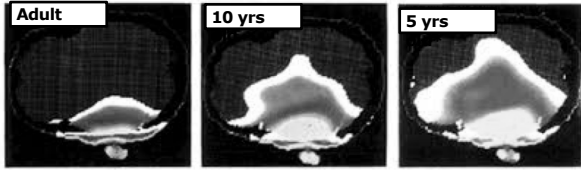
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Ex.1. Light traveling in air encounters the water; another medium.

Ex. 3. Wave encounters skin tissue, muscle mass and bones

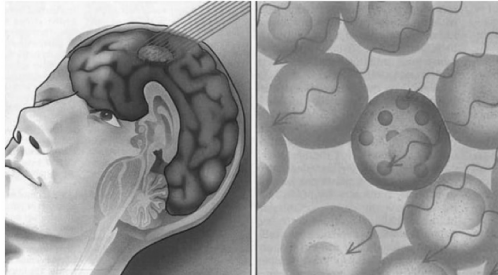
- Cellphone radiation



*Kid's skull is not yet developed

Ex. 4. Cancer treatment

Darren Quick | April 3, 2013

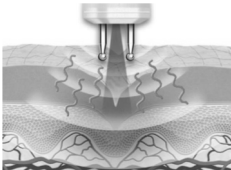


Boron neutron capture therapy can kill tumors without harming healthy neighboring tissue.

Wave incidence

- For many applications, [such as fiber optics, line power transmission], it's necessary to know what happens to a wave when it meets a different medium.

- How much is *transmitted*?
- How much is *reflected* back?



We will look at...

I. Normal incidence

Wave arrives at 0° from normal

- Standing waves



II. Oblique incidence

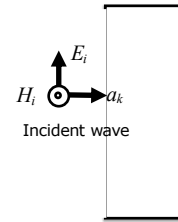
Wave arrives at another angle

- Snell's Law and Critical angle
- Parallel or Perpendicular polarization
- Brewster angle

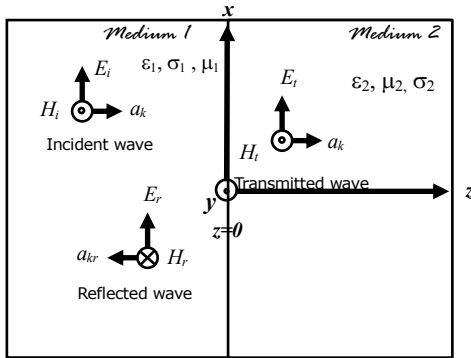


Vital Definition

- Plane between media-interface
- Plane of incidence- (what you draw)



Reflection at Normal Incidence

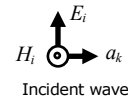
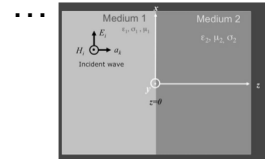


Now in terms of equations

- Incident wave

$$\vec{E}_{is}(z) = E_{io} e^{-\gamma_1 z} \hat{x}$$

$$\vec{H}_{is}(z) = H_{io} e^{-\gamma_1 z} \hat{y} = \frac{E_{io}}{\eta_1} e^{-\gamma_1 z} \hat{y}$$

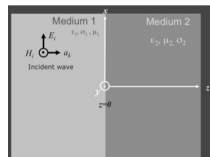
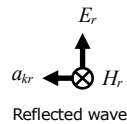


Reflected wave

- It's traveling along $-z$ axis

$$\vec{E}_{rs}(z) = E_{ro} e^{\gamma_1 z} \hat{x}$$

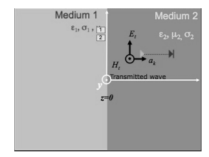
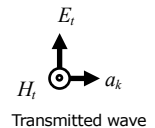
$$\vec{H}_{rs}(z) = H_{ro} e^{\gamma_1 z} (-\hat{y}) = -\frac{E_{ro}}{\eta_1} e^{\gamma_1 z} \hat{y}$$



Transmitted wave

$$\vec{E}_{ts}(z) = E_{to} e^{-\gamma_2 z} \hat{x}$$

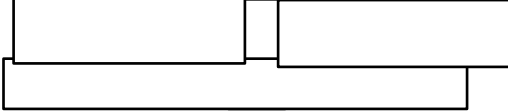
$$\vec{H}_{ts}(z) = H_{to} e^{-\gamma_2 z} \hat{y} = \frac{E_{to}}{\eta_2} e^{-\gamma_2 z} \hat{y}$$



Sum to get TOTAL field



- At medium 1 and medium 2



- Tangential components must be continuous at the interface

$$\vec{E}_i(0) + \vec{E}_r(0) = \vec{E}_t(0)$$

$$\vec{H}_i(0) + \vec{H}_r(0) = \vec{H}_t(0)$$

Look at E fields

$$\vec{E}_i(0) + \vec{E}_r(0) = \vec{E}_t(0)$$

$$E_{io} e^{-\gamma_1 z} \hat{x} + E_{ro} e^{\gamma_1 z} \hat{x} = E_{to} e^{-\gamma_2 z} \hat{x}$$

$$z = 0$$

$$E_{io} + E_{ro} = E_{to}$$

$$\frac{E_{io}}{E_{io}} + \frac{E_{ro}}{E_{io}} = \frac{E_{to}}{E_{io}}$$

Define

- Reflection coefficient, Γ

$$\Gamma = \frac{E_{ro}}{E_{io}} = \frac{\eta_2 - \eta_1}{\eta_2 + \eta_1}$$

- Transmission coefficient, τ

$$\tau = \frac{E_{to}}{E_{io}} = \frac{2\eta_2}{\eta_2 + \eta_1}$$

Note:

- $1 + \Gamma = \tau$
- Both are dimensionless and may be complex
- $0 \leq |\Gamma| \leq 1$

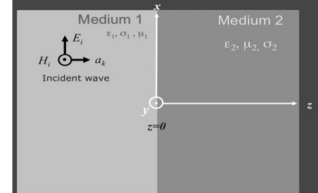
Example: PE 10.8

A 5GHz uniform plane wave $E_{is} = 10e^{-j\beta z} \hat{a}_x$ in free space is incident normally on a large plane, lossless dielectric slab ($z > 0$) having $\epsilon = 4\epsilon_0$ and $\mu = \mu_0$.

Find:

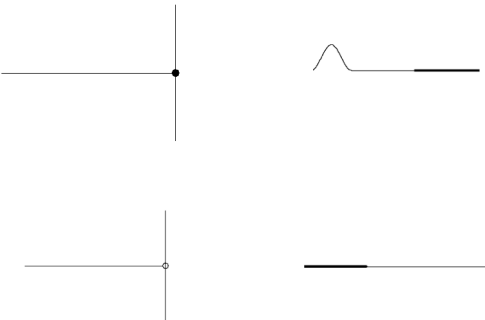
- the reflected wave E_{rs} and
- the transmitted wave E_{ts} .

Answer:



SEE
<http://www.acs.psu.edu/drussell/Demos/reflect/reflect.html>

<http://www.acs.psu.edu/drussell/Demos/reflect/reflect.html>



Today we'll see 3 Cases

and the Standing waves form on each case

- Medium 2 is perfect Conductor
- Medium 2 is perfect Dielectric
 - $\eta_2 > \eta_1$
- Medium 2 is perfect Dielectric
 - $\eta_1 > \eta_2$

Case 1: Conductor(2)

- Medium 1: perfect dielectric, $\sigma_1=0$
 - Medium 2: perfect conductor, $\sigma_2=\infty$
- Find intrinsic impedance
Coef. Of reflex & transmission
& E1 field

$$\eta_2 = 0, \quad \Gamma = -1, \tau = 0$$

$$E_{i0} = -2jE_{i0} \sin \beta_1 z \hat{x} \text{ (phasor)}$$

$$E_1(z,t) = 2E_{i0} \sin \beta_1 z \sin \omega t \hat{x}$$

<http://www.phy.ntnu.edu.tw/java/waveSuperposition/waveSuperposition.html>

The EM field forms a Standing Wave on medium 1

$$E_1 = 2E_{i0} \sin \beta_1 z \sin \omega t \hat{x}$$

Minima @ $-\beta_1 z = 0, \pi, 2\pi$
Maxima @ $-\beta_1 z = \frac{\pi}{2}, \frac{3\pi}{2}, \frac{5\pi}{2}$
 $z_{\max} = -\frac{n\pi}{\beta_1} = -\frac{n\lambda_1}{4} \quad n = 1, 3, 5$

Case 2: Dielectric 2 $\eta_2 > \eta_1$

- ◆ Medium 1: perfect dielectric $\sigma_1=0$
- ◆ Medium 2: perfect dielectric $\sigma_2=0, \eta_2 > \eta_1$

If $\eta_2 > \eta_1$,
 $\Gamma > 0$,
 τ and Γ are real.

$$E_{1s} = E_{is} + E_{rs}$$

$$= E_{oi}(e^{-j\beta_1 z} + \Gamma e^{+j\beta_1 z})$$

$$= E_{oi}e^{-j\beta_1 z}(1 + \Gamma e^{+2j\beta_1 z})$$

$$-2\beta_1 z_{\max} = 0, 2\pi, 4\pi, 6\pi, \dots$$

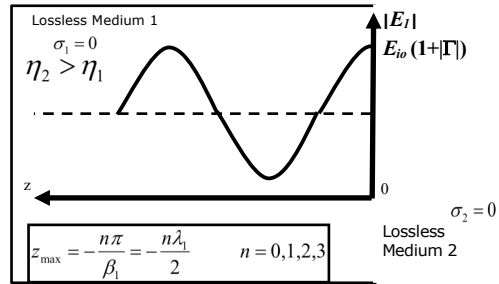
or $-\beta_1 z_{\max} = 0, \pi, 2\pi, 3\pi, \dots$

$$z_{\max} = -\frac{n\pi}{\beta_1} = -\frac{n\lambda_1}{2} \quad n = 0, 1, 2, 3$$

$$z_{\min} = -\frac{(2n+1)\pi}{2\beta_1} = -\frac{(2n+1)\lambda_1}{4} \quad n = 0, 1, 2, 3$$

Standing waves due to reflection

$$E_1 = E_i + E_r = E_{oi}(e^{-j\beta_1 z} + \Gamma e^{+j\beta_1 z}) = E_{oi}e^{-j\beta_1 z}(1 + \Gamma e^{+2j\beta_1 z})$$



At every half-wavelength, everything repeats!

Case 3: Dielectric 2 $\eta_1 > \eta_2$

- ◆ Medium 1 = perfect dielectric $\sigma_1=0$
- Medium 2 = perfect dielectric $\sigma_2=0, \eta_1 > \eta_2$

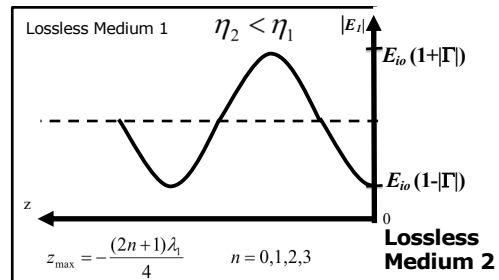
If $\eta_2 < \eta_1$,
 $\Gamma < 0$, τ and Γ are real.

$$z_{\max} = -\frac{(2n+1)\pi}{2\beta_1} = -\frac{(2n+1)\lambda_1}{4} \quad n = 1, 2, 3$$

$$z_{\min} = -\frac{n\pi}{\beta_1} = -\frac{n\lambda_1}{2} \quad n = 0, 1, 2, 3$$

Standing waves due to reflection

$$E_1 = E_i + E_r = E_{oi}(e^{-j\beta_1 z} + \Gamma e^{+j\beta_1 z}) = E_{oi}e^{-j\beta_1 z}(1 + \Gamma e^{+2j\beta_1 z})$$



At every half-wavelength, all em properties repeat

Standing Wave Ratio, s

- Measures the amount of reflections, the more reflections, the larger the standing wave that is formed.
- The ratio of $|E_1|_{max}$ to $|E_1|_{min}$

$$s = \frac{|E_1|_{max}}{|E_1|_{min}} = \frac{|H_1|_{max}}{|H_1|_{min}} = \frac{1 + |\Gamma|}{1 - |\Gamma|}$$

Ideally $s=1$ (0 dB)
No reflections

or

$$|\Gamma| = \frac{s - 1}{s + 1}$$

Ex. Given a wave travelling in F.S. and hitting dielectric

- Lossless dielectric at 10MHz has $\epsilon_r=16$, nonmagnetic
- Find SWR on air

Power Flow in Medium 1

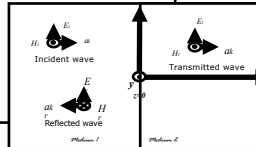
- The net average power density flowing in lossless medium 1

$$P_{ave1}(z) = \frac{1}{2} \text{Re}[E_1 \times H_1^*]$$

$$= \frac{1}{2} \text{Re} \left[\hat{x} E_{io} (e^{-j\beta_1 z} + \Gamma e^{j\beta_1 z}) \times \hat{y} \frac{E_{io}^*}{\eta_1} (e^{j\beta_1 z} + \Gamma^* e^{-j\beta_1 z}) \right]$$

$$= \hat{z} \frac{|E_{io}|^2}{2\eta_1} (1 - |\Gamma|^2)$$

$$= P_{ave}^i + P_{ave}^r$$



Power Flow in Transmitted wave

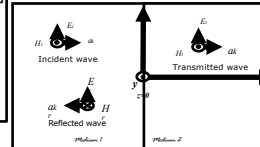
- The net average power density flowing in lossless medium 2

$$P_{ave2}(z) = P_t = \frac{1}{2} \text{Re}[E_2 \times H_2^*]$$

$$= \frac{1}{2} \text{Re} \left[\hat{x} \tau E_{io} e^{-j\beta_2 z} \times \hat{y} \tau^* \frac{E_{io}^*}{\eta_2} e^{j\beta_2 z} \right]$$

$$P_t = \hat{z} |\tau|^2 \frac{|E_{io}|^2}{2\eta_2}$$

where $\hat{a}_k = \hat{z}$



Power in Lossy Media

$$P_{ave1}(z) = \hat{z} \frac{|E_o^i|^2}{2\eta_{cl}} (e^{-2\alpha_1 z} - |\Gamma|^2 e^{2\alpha_1 z})$$

$$P_{ave2}(z) = \hat{z} |\tau|^2 \frac{|E_o^i|^2}{2} e^{-2\alpha_2 z} \text{Re} \left(\frac{1}{\eta_{c2}^*} \right)$$

where

$$\Gamma = \frac{\sqrt{\epsilon_{c1}} - \sqrt{\epsilon_{c2}}}{\sqrt{\epsilon_{c1}} + \sqrt{\epsilon_{c2}}} \quad \text{and} \quad \epsilon_{c2} = \epsilon_2 - j \frac{\sigma_2}{\omega_2}$$

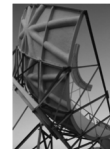
Ex. Antenna Radome



Antenna with radome

A 10GHz aircraft radar uses a narrow-beam scanning antenna mounted on a gimbal behind a dielectric radome.

- Even though the radome shape is far from planar, it is approximately planar over the narrow extent of the radar beam.
- If the radome material is a lossless dielectric with $\mu_r=1$ and $\epsilon_r=9$, choose its thickness d such that the radome appears transparent to the radar beam.
- Mechanical integrity requires d to be greater than 2.3 cm.



Antenna with no radome

Answer:
 $\lambda/2 = .5\text{cm}, d = 2.5\text{cm}$