




Electromagnetism

INEL 4151



Sandra Cruz-Pol, Ph. D.
ECE UPRM
Mayaguez, PR



In summary

- Stationary Charges
 - Q
- Steady currents
 - I
- Time-varying currents
 - $I(t)$
- Electrostatic fields\
 - E
- Magnetostatic fields
 - H
- Electromagnetic (waves!)
 - $E(t)$ & $H(t)$

Cruz-Pol, Electromagnetics
UPRM

Outline

- Faraday's Law & Origin of emag
- Transformer and Motional EMF
- Displacement Current & Maxwell Equations
- Review: Phasors and Time Harmonic fields

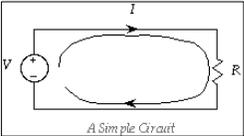


9.2

Faraday's Law

Electricity => Magnetism

- In 1820 Oersted discovered that a steady current produces a magnetic field while teaching a physics class.



A Simple Circuit



This is what Oersted discovered accidentally:

$$\oint_L \vec{H} \cdot d\vec{l} = \int_s \vec{J} \cdot d\vec{S}$$

Cruz-Pol, Electromagnetics
UPRM

Would magnetism would produce electricity?

- Eleven years later, and at the same time, (Mike) Faraday in London & (Joe) Henry in New York discovered that a time-varying magnetic field would produce an electric current!



$$V_{emf} = -N \frac{d\Psi}{dt}$$

$$\oint_L \vec{E} \cdot d\vec{l} = -N \int_s \frac{\partial}{\partial t} \vec{B} \cdot d\vec{S}$$

Cruz-Pol, Electromagnetics
UPRM

Len's Law = (-)

➤ If $N=1$ (1 loop)

$$V_{emf} = -\frac{d\Psi}{dt}$$

➤ The time change

$$V_{emf} = \oint_L \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \int_s \vec{B} \cdot d\vec{S}$$

can refer to B or S

ctromagnetics
UPRM

Electromagnetics was born!

➤ This is Faraday's Law - the principle of motors, hydro-electric generators and transformers operation.

Faraday's Law $\oint_L \vec{E} \cdot d\vec{l} = -\int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S}$

Ampere's Law $\oint_L \vec{H} \cdot d\vec{l} = \int_s \vec{J} \cdot d\vec{S}$

**Mention some examples of em waves*

Cruz-Pol, Electromagnetics
UPRM

Faraday's Law

➤ For $N=1$ and $B=0$

A Simple Circuit

$$V_{emf} = -N \frac{d\Psi}{dt}$$

$$V_{emf} = \oint_L \vec{E} \cdot d\vec{l} = IR$$

Cruz-Pol, Electromagnetics
UPRM

Example PE 9.3 A magnetic core of uniform cross-section 4 cm^2 is connected to a 120V, 60Hz generator. Calculate the induced emf V_2 in the secondary coil. $N_1=500, N_2=300$

➤ Use Faraday's Law

$$V_1 = -N_1 \frac{d\Psi}{dt}$$

$$V_2 = -N_2 \frac{d\Psi}{dt}$$

$$V_2 = N_2 \frac{V_1}{N_1}$$

Transformer Core

Answer: $72 \cos(120\pi t) \text{ V}$

9.3
Transformer & Motional EMF

Two cases of $V_{emf} = \oint_L \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \int_s \vec{B} \cdot d\vec{S}$

➤ B changes

$$\oint_L \vec{E} \cdot d\vec{l} = -\int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S}$$

Stoke's theorem

$$\int_s (\nabla \times \vec{E}) \cdot d\vec{S} = -\int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

➤ S (area) changes

$$\oint_L \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \int_s \vec{B} \cdot d\vec{S}$$

$$\vec{E}_m = \vec{F}_m / Q = \vec{u} \times \vec{B}$$

$$\int_s (\nabla \times \vec{E}) \cdot d\vec{S} = \int_L (\vec{u} \times \vec{B}) \cdot d\vec{l}$$

$$\nabla \times \vec{E} = \nabla \times \vec{u} \times \vec{B}$$

Cruz-Pol, Electromagnetics
UPRM

3 cases:

- **Stationary loop in time-varying B field**

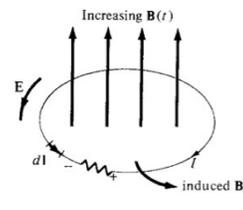
$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$
- **Moving loop in static B field**

$$\nabla \times \vec{E} = \nabla \times \vec{u} \times \vec{B}$$
- **Moving loop in time-varying B field**

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} + \nabla \times \vec{u} \times \vec{B}$$

Cruz-Pol, Electromagnetics UPRM

1. Stationary loop in a time-varying B field

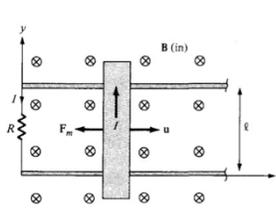


$$V_{emf} = \oint_L \vec{E} \cdot d\vec{l} = -\int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

Cruz-Pol, Electromagnetics UPRM

2. Time-varying loop area in a static B field



$$\vec{E}_m = \frac{\vec{F}_m}{Q} = \vec{u} \times \vec{B}$$

$$V_{emf} = \oint_L \vec{E}_m \cdot d\vec{l} = \oint_L (\vec{u} \times \vec{B}) \cdot d\vec{l}$$

$$\nabla \times \vec{E}_m = \nabla \times (\vec{u} \times \vec{B})$$

Cruz-Pol, Electromagnetics UPRM

3. A t-varying loop area in a t-varying B field

$$V_{emf} = \oint_L \vec{E} \cdot d\vec{l} = -\int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} + \oint_L (\vec{u} \times \vec{B}) \cdot d\vec{l}$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} + \nabla \times (\vec{u} \times \vec{B})$$

Cruz-Pol, Electromagnetics UPRM

PE 9.1

A conducting bar can slide freely over two conducting rails as shown in Figure 9.6. Calculate the induced voltage in the bar

- If the bar is stationed at $y = 8$ cm and $\vec{B} = 4 \cos(10^6 t) \vec{a}_z$ mWb/m²
- If the bar slides at a velocity $\vec{u} = 20\vec{a}_x$ m/s and $\vec{B} = 4\vec{a}_z$ mWb/m²
- If the bar slides at a velocity $\vec{u} = 20\vec{a}_x$ m/s and $\vec{B} = 4 \cos(10^6 t - y) \vec{a}_z$ mWb/m²

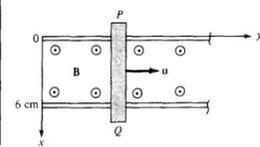


Figure 9.6 For Example 9.1.

$$V_{emf} = -\int_s \frac{\partial \vec{B}}{\partial t} \cdot d\vec{S} = 4 \cdot 10^{-6} \sin(10^6 t) (0.08)(0.06) = 19.2 \sin(10^6 t) \mu V$$

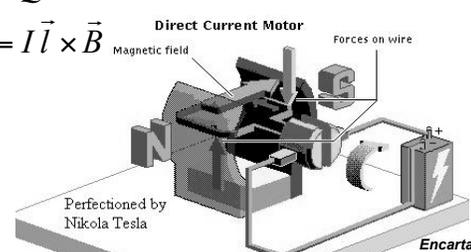
$$V_{emf} = \oint_L \vec{E} \cdot d\vec{l} = \oint_L \vec{u} \times \vec{B} \cdot d\vec{l}$$

Cruz-Pol, Electromagnetics UPRM

Moving loop in static B field

- When a conducting loop is moving inside a magnet (static B field), the force on a charge is:

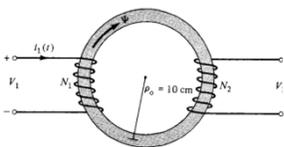
$$\vec{F} = Q\vec{u} \times \vec{B}$$

$$\vec{F} = I\vec{l} \times \vec{B}$$


Cruz-Pol, Electromagnetics UPRM

Transformer Example

The magnetic circuit of Figure 9.8 has a uniform cross section of 10^{-3} m^2 . If the circuit is energized by a current $i_1(t) = 3 \sin 100\pi t \text{ A}$ in the coil of $N_1 = 200$ turns, find the emf induced in the coil of $N_2 = 100$ turns. Assume that $\mu = 500 \mu_0$.



$$\Psi = \frac{\mathcal{F}}{\mathcal{R}} = \frac{N_1 i_1}{\ell / \mu S} = \frac{N_1 i_1 \mu S}{2\pi \rho_o}$$

➤ Find reluctance and use Faraday's Law

$$V_2 = -N_2 \frac{d\Psi}{dt} = -\frac{N_1 N_2 \mu S}{2\pi \rho_o} \frac{di_1}{dt}$$

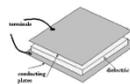
$$V_{emf} = -100(200)(3\omega \cos \omega t) \frac{500 \mu_0 10^{-3}}{2\pi(0.10)} = -6\pi \cos 100\pi t$$

9.4

Displacement Current, J_d

Maxwell noticed something was missing...

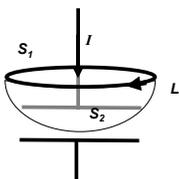
➤ And added J_d , the displacement current



$$\oint_L H \cdot dl = \int_{S_1} J \cdot dS = I_{enc} = I$$

$$\oint_L H \cdot dl = \int_{S_2} J \cdot dS = 0$$

$$\oint_L H \cdot dl = \int_{S_2} J_d \cdot dS = \frac{d}{dt} \int_{S_2} D \cdot dS = \frac{dQ}{dt} = I$$



At low frequencies $J \gg J_d$, but at radio frequencies both terms are comparable in magnitude.

9.4

Maxwell's Equation in Final Form

Summary of Terms

- E = electric field intensity [V/m]
- D = electric field density [C/m^2]
- H = magnetic field intensity, [A/m]
- B = magnetic field density, [Teslas]
- J = current density [A/m^2]

Cruz-Pol, Electromagnetics UPRM

Maxwell Equations in General Form

Differential form	Integral Form	
$\nabla \cdot D = \rho_v$	$\oint_S D \cdot dS = \int_V \rho_v dv$	Gauss's Law for E field.
$\nabla \cdot B = 0$	$\oint_S B \cdot dS = 0$	Gauss's Law for H field. Nonexistence of monopole
$\nabla \times E = -\frac{\partial B}{\partial t}$	$\oint_L E \cdot dl = -\frac{\partial}{\partial t} \int_S B \cdot dS$	Faraday's Law
$\nabla \times H = J + \frac{\partial D}{\partial t}$	$\oint_L H \cdot dl = \int_S \left(J + \frac{\partial D}{\partial t} \right) \cdot dS$	Ampere's Circuit Law

Cruz-Pol, Electromagnetics UPRM

Maxwell's Eqs.

- Also the equation of continuity $\nabla \cdot \mathbf{J} = -\frac{\partial \rho_v}{\partial t}$
- Maxwell added the term $\frac{\partial D}{\partial t}$ to Ampere's Law so that it not only works for **static** conditions but also for **time-varying** situations.
 - This added term is called the displacement current density, while \mathbf{J} is the conduction current.

Cruz-Pol, Electromagnetics
UPRM

Relations & B.C.

$$\mathbf{D} = \epsilon \mathbf{E} = \epsilon_0 \mathbf{E} + \mathbf{P}$$

$$\mathbf{B} = \mu \mathbf{H} = \mu_0 (\mathbf{H} + \mathbf{M})$$

$$\mathbf{J} = \sigma \mathbf{E} + \rho_v \mathbf{u}$$

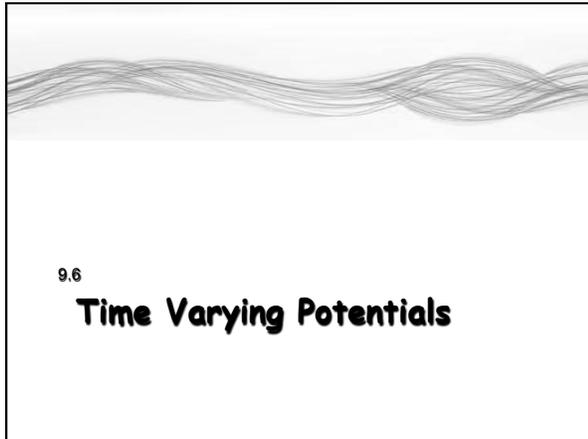
$$E_{1t} = E_{2t} \quad \text{or} \quad (\mathbf{E}_1 - \mathbf{E}_2) \times \mathbf{a}_{n12} = 0$$

$$H_{1t} - H_{2t} = K \quad \text{or} \quad (\mathbf{H}_1 - \mathbf{H}_2) \times \mathbf{a}_{n12} = \mathbf{K}$$

$$D_{1n} - D_{2n} = \rho_s \quad \text{or} \quad (\mathbf{D}_1 - \mathbf{D}_2) \cdot \mathbf{a}_{n12} = \rho_s$$

$$B_{1n} - B_{2n} = 0 \quad \text{or} \quad (\mathbf{B}_2 - \mathbf{B}_1) \cdot \mathbf{a}_{n12} = 0$$

UPRM



We had defined

- **Electric & Magnetic potentials:**

$$V = \int_v \frac{\rho_v dv}{4\pi\epsilon R} \quad \mathbf{A} = \int_v \frac{\mu \mathbf{J} dv}{4\pi R}$$

- **Related to B as:**

$$\mathbf{B} = \nabla \times \mathbf{A}$$

- Substituting into Faraday's law:

$$\nabla \times \mathbf{E} = -\frac{\partial}{\partial t} (\nabla \times \mathbf{A})$$

$$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t}$$

Cruz-Pol, Electromagnetics
UPRM

Electric & Magnetic potentials:

- If we take the divergence of \mathbf{E} :

$$\nabla \cdot \mathbf{E} = \frac{\rho_v}{\epsilon} = -\nabla^2 V - \frac{\partial}{\partial t} (\nabla \cdot \mathbf{A})$$

- Or

$$\nabla^2 V + \frac{\partial}{\partial t} (\nabla \cdot \mathbf{A}) = -\frac{\rho_v}{\epsilon}$$

- Taking the curl of: $\mathbf{B} = \nabla \times \mathbf{A}$ & add Ampere's we get

$$\begin{aligned} \nabla \times \nabla \times \mathbf{A} &= \mu \mathbf{J} + \epsilon \mu \frac{\partial}{\partial t} \left(-\nabla V - \frac{\partial \mathbf{A}}{\partial t} \right) \\ &= \mu \mathbf{J} - \mu \epsilon \nabla \left(\frac{\partial V}{\partial t} \right) - \mu \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} \end{aligned}$$

Cruz-Pol, Electromagnetics
UPRM

Electric & Magnetic potentials:

- If we apply this vector identity

$$\nabla \times \nabla \times \mathbf{A} = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$$

- We end up with

$$\nabla^2 \mathbf{A} - \nabla(\nabla \cdot \mathbf{A}) = -\mu \mathbf{J} + \mu \epsilon \nabla \left(\frac{\partial V}{\partial t} \right) + \mu \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2}$$

Cruz-Pol, Electromagnetics
UPRM

Electric & Magnetic potentials:

➤ We use the Lorentz condition:

$$\nabla \cdot \mathbf{A} = -\mu\epsilon \frac{\partial V}{\partial t}$$

To get:

$$\nabla^2 \mathbf{A} - \mu\epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} = -\mu\mathbf{J}$$

and:

$$\nabla^2 V - \mu\epsilon \frac{\partial^2 V}{\partial t^2} = -\frac{\rho_v}{\epsilon}$$

Which are both wave equations.

Cruz-Pol, Electromagnetics
UPRM



9.7

Time Harmonic Fields Phasors Review

Time Harmonic Fields

➤ Definition: is a field that varies periodically with time.

- Ex. Sinusoid

➤ Let's review Phasors!

Cruz-Pol, Electromagnetics
UPRM

Phasors & complex #'s

Working with harmonic fields is easier, but requires knowledge of phasor, let's review

➤ **complex numbers** and

➤ **phasors**

Cruz-Pol, Electromagnetics
UPRM

COMPLEX NUMBERS:

➤ Given a complex number z

$$z = x + jy = re^{j\phi} = r\angle\phi = r \cos\phi + jr \sin\phi$$

where $r = |z| = \sqrt{x^2 + y^2}$ is the magnitude

$$\phi = \tan^{-1} \frac{y}{x}$$

is the angle

Cruz-Pol, Electromagnetics
UPRM

Review:

examples :

➤ Addition, $1/j =$

➤ Subtraction, $1/\sqrt{j} =$

➤ Multiplication, $e^{j45^\circ} / j =$

➤ Division, $e^{j45^\circ} - \sqrt{j} =$

➤ Square Root, $3e^{j90^\circ} + j =$

➤ Complex Conjugate $\sqrt{2}e^{+j45^\circ} =$

$$\sqrt{2}e^{+j45^\circ} + 10e^{+j90^\circ} =$$

Cruz-Pol, Electromagnetics
UPRM

For a Time-varying phase

$$\phi = \omega t + \theta$$

Real and imaginary parts are:

$$\text{Re}\{re^{j\phi}\} = r \cos(\omega t + \theta)$$

$$\text{Im}\{re^{j\phi}\} = r \sin(\omega t + \theta)$$

Cruz-Pol, Electromagnetics
UPRM

PHASORS

➤ For a sinusoidal current $I(t) = I_o \cos(\omega t + \theta)$ equals the real part of $I_o e^{j\theta} e^{j\omega t}$

➤ The complex term $I_o e^{j\theta}$ which results from dropping the time factor $e^{j\omega t}$ is called the phasor current, denoted by I_s (s comes from sinusoidal)

Cruz-Pol, Electromagnetics
UPRM

To change back to time domain

The phasor is

1. multiplied by the time factor, $e^{j\omega t}$,
2. and taken the real part.

$$A = \text{Re}\{A_s e^{j\omega t}\}$$

Cruz-Pol, Electromagnetics
UPRM

Advantages of phasors

➤ **Time derivative** in time is equivalent to multiplying its phasor by $j\omega$

$$\frac{\partial A}{\partial t} \rightarrow j\omega A_s$$

➤ **Time integral** is equivalent to dividing by the same term.

$$\int A dt \rightarrow \frac{A_s}{j\omega}$$

Cruz-Pol, Electromagnetics
UPRM



9.7

Time Harmonic Fields

Time-Harmonic fields (sines and cosines)

➤ The wave equation can be derived from Maxwell equations, indicating that the changes in the fields behave as a wave, called an electromagnetic wave or field.

➤ Since any periodic wave can be represented as a sum of sines and cosines (using Fourier), then we can deal only with harmonic fields to simplify the equations.

Cruz-Pol, Electromagnetics
UPRM

Maxwell Equations for Harmonic fields (phasors)	
Differential form*	
$\nabla \cdot \epsilon E = \rho_v$	Gauss' s Law for E field.
$\nabla \cdot \mu H = 0$	Gauss' s Law for H field. No monopole
$\nabla \times E = -j\omega\mu H$	Faraday' s Law
$\nabla \times H = \sigma E + j\omega\epsilon E$	Ampere' s Circuit Law
* (substituting $D = \epsilon E$ and $H = \mu B$)	

