

Electrostatic fields



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Some applications

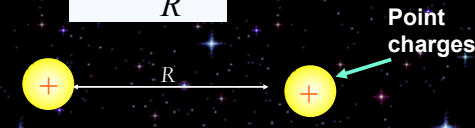
- Power transmission, X rays, lightning protection
- Solid-state Electronics: resistors, capacitors, FET
- Computer peripherals: touch pads, LCD, CRT
- Medicine: electrocardiograms, electroencephalograms, monitoring eye activity
- Agriculture: seed sorting, moisture content monitoring, spinning cotton, ...
- Art: spray painting
- ...

We will study Electric charges:

- **Coulomb's Law** - $F = \frac{kQ_1Q_2}{R^2}$
 - Use when charge distribution is known
- **Gauss' s Law** - $\Psi = \oint_S \vec{D} \cdot d\vec{S} = Q_{enc}$
 - Use when charge distribution is symmetrical
- **Electric Potential**
(uses scalar, not vectors) $V(r) = -\int_{\infty}^r \vec{E} \cdot d\vec{l}$
 - Use when potential V is known

Coulomb's Law (1785)

- Force one charge exerts on another

$$F = \frac{kQ_1Q_2}{R^2}$$


where $k = 9 \times 10^9$
or $k = 1/4\pi\epsilon_0$
 $\epsilon_0 = 8.85 \times 10^{-12}$

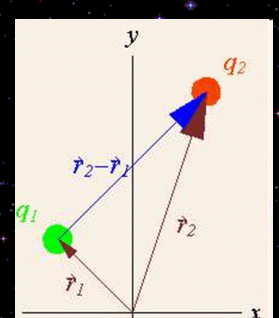
*Superposition applies

Force with direction

Force that Q_1 exerts on Q_2

$$\vec{F}_{12} = \frac{Q_1Q_2}{4\pi\epsilon_0 R^2} \hat{a}_{12}$$

Note: Observation point goes First!

$$\hat{a}_{12} = \frac{\vec{r}_2 - \vec{r}_1}{|\vec{r}_2 - \vec{r}_1|}$$


Example

Example: Point charges 5nC and -2nC are located at $r_1=(2,0,4)$ and $r_2=(-3,0,5)$, respectively.

a) Find the force on a 1nC point charge, Q_3 , located at $(1,-3,7)$

Apply superposition:

$$\vec{F}_x = \frac{1}{4\pi\epsilon_0} \left[\frac{Q_1Q_3(\vec{r}_x - \vec{r}_1)}{|\vec{r}_x - \vec{r}_1|^3} + \frac{Q_2Q_3(\vec{r}_x - \vec{r}_2)}{|\vec{r}_x - \vec{r}_2|^3} \right]$$

Electric field intensity

- Is the force per unit charge when placed in the E field

$$E = \frac{F}{Q}$$

$$E = \frac{Q}{4\pi\epsilon_0 R^2} \hat{a}_R$$

Example: Same point charges 5nC and -2nC are located at (2,0,4) and (-3,0,5), respectively.

b) Find the E field at $r_s = (1, -3, 7)$.

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \left[\frac{Q_1(r_s - r_1)}{|r_s - r_1|^3} + \frac{Q_2(r_s - r_2)}{|r_s - r_2|^3} \right] = (-1.0, -1.29, 1.4) \text{ V/m}$$

If we have many charges

| | |
|------------------------------------|------------------|
| Line charge density, ρ_L | C/m |
| Surface charge density ρ_S | C/m ² |
| Volume charge density ρ_V | C/m ³ |

$E = \frac{F}{Q}$
 $E = \frac{Q}{4\pi\epsilon_0 R^2} \hat{a}_R$

$Q = \int_L \rho_L dl$

$Q = \int_S \rho_S dS$

$Q = \int_V \rho_V dv$

The total E-field intensity is

$$E_{\text{point charge}} = \frac{Q}{4\pi\epsilon_0 R^2} \hat{a}_R$$

$$E = \int \frac{(\rho_L dl)}{4\pi\epsilon_0 R^2} \hat{a}_R$$

$$E = \int \frac{(\rho_S dS)}{4\pi\epsilon_0 R^2} \hat{a}_R$$

$$E = \int \frac{(\rho_V dv)}{4\pi\epsilon_0 R^2} \hat{a}_R$$

More Charge distributions

Find E from

- Point charge (we just saw this one)
- Line charge
- Surface charge
- Volume charge

Results Preview

| | | |
|---------------|---|--|
| Line charge | $\vec{E} = \frac{\rho_L}{2\pi\epsilon_0 \rho} \hat{a}_\rho$ | We will derive these 3 cases 1. Using Coulomb 2. Using Gauss |
| Sheet charge | $\vec{E} = \frac{\rho_S}{2\epsilon_0} \hat{a}_n$ | |
| Volume Charge | $\vec{E} = \frac{Q}{4\pi\epsilon_0 r^2} \hat{a}_r$ | |

Find E from LINE charge

- Line charge w/ uniform charge density, ρ_L

*use cylindrical coordinates

$$Q = \int_A^B \rho_L dl \quad dl = dz'$$

$$E = \int \frac{\rho_L dz'}{4\pi\epsilon_0 R^2} \hat{a}_R$$

$$z' = OT - \rho \tan \alpha$$

$$R = \rho \sec \alpha$$

$$\vec{R} = R \cos \alpha \hat{a}_\rho + R \sin \alpha \hat{a}_z$$

$$\hat{a}_R = \frac{\vec{R}}{R} = \cos \alpha \hat{a}_\rho + \sin \alpha \hat{a}_z$$

Defining angles α_1 and α_2

α_1 = imaginary perpendicular line with the **back**
 α_2 = imaginary perpendicular line with the **front**

LINE charge

Substituting in: $E = \int \frac{\rho_L dz'}{4\pi\epsilon_0 R^2} \hat{a}_R$

$z' = OT - \rho \tan \alpha$
 $dz' = [0 - \rho \sec^2 \alpha] d\alpha$
 $R = \rho \sec \alpha$
 $\hat{a}_R = \cos \alpha \hat{a}_\rho + \sin \alpha \hat{a}_z$

$$\vec{E} = \int_{\alpha_1}^{\alpha_2} \frac{\rho_L [-\rho \sec^2 \alpha] d\alpha}{4\pi\epsilon_0 \rho^2 \sec^2 \alpha} [\cos \alpha \hat{a}_\rho + \sin \alpha \hat{a}_z]$$

finite Line Charge:

$$\vec{E} = \frac{\rho_L}{4\pi\epsilon_0 \rho} [-(\sin \alpha_2 - \sin \alpha_1) \hat{a}_\rho + (\cos \alpha_2 - \cos \alpha_1) \hat{a}_z]$$

infinite Line Charge ($\alpha_{1,2} = \mp 90^\circ$)

$$E = \frac{\rho_L}{2\pi\epsilon_0 \rho} \hat{a}_\rho$$

More Charge distributions

- Point charge
- Line charge
- **Surface charge**
- Volume charge

Find E from Surface charge

Sheet of charge w/ uniform density ρ_s

$dQ = \rho_s dS$ $dE = \frac{\rho_s dS}{4\pi\epsilon_0 R^2} \hat{a}_R$

Element of area is: $dS = \rho d\phi d\rho$

Observation point is at z-axis:
 $R = \rho(-\hat{a}_\rho) + h\hat{a}_z$

$$\hat{a}_R = \frac{\vec{R}}{R}$$

$$dE = \frac{\rho_s \rho d\phi d\rho [-\rho \hat{a}_\rho + h\hat{a}_z]}{4\pi\epsilon_0 [\rho^2 + h^2]^{3/2}}$$

SURFACE charge

Due to SYMMETRY the ρ component cancels out.

$$E_z = \frac{\rho_s}{4\pi\epsilon_0} \int_{\phi=0}^{2\pi} d\phi \int_{\rho=0}^{\infty} \frac{h\rho d\rho}{[\rho^2 + h^2]^{3/2}}$$

$$E_z = \frac{\rho_s}{4\pi\epsilon_0} 2\pi \left[\frac{-h}{\sqrt{\rho^2 + h^2}} \right]_0^{\infty}$$

infinite Surface Charge:

$$E = \frac{\rho_s}{2\epsilon_0} \hat{a}_n$$

More Charge distributions

- Point charge
- Line charge
- Surface charge
- **Volume charge**

Find E from Volume charge

- Sphere of charge w/ uniform density, ρ_v

$$dQ = \rho_v dv$$

$$dE = \frac{\rho_v dv}{4\pi\epsilon_0 R^2} \hat{a}_R$$

$$dv = r'^2 \sin \theta' d\theta' d\phi' dr'$$

Law of cosines:

$$R^2 = z^2 + r'^2 - 2zr' \cos \theta'$$

$$r'^2 = z^2 + R^2 - 2zR \cos \alpha$$

Differentiating (Eq. *)

$$\sin \theta' d\theta' = \frac{RdR}{zr'}$$

Due to symmetry only $dE_z = dE \cos \alpha$ survives.

Find E from Volume charge

- Substituting...

$$dE_z = \frac{\rho_v dv}{4\pi\epsilon_0 R^2} \cos \alpha \hat{a}_z$$

$$r'^2 = z^2 + R^2 - 2zR \cos \alpha$$

$$dv = r'^2 \sin \theta' d\theta' d\phi' dr'$$

$$\sin \theta' d\theta' = \frac{RdR}{zr'}$$

$$E_z = \frac{\rho_v}{4\pi\epsilon_0} \int_{\phi=0}^{2\pi} \int_{\theta=0}^{\pi} \int_{r=0}^R \frac{1}{4\pi\epsilon_0 z^2} \left(\frac{4}{3} \pi a^3 \rho_v \right) \hat{a}_z = \frac{Q}{4\pi\epsilon_0 r^2} \hat{a}_z - \frac{1}{R^2}$$

De donde salen los límites de R?

P.E. 4.5

- A square plate at plane $z=0$ and $x \leq \pm 2$, $y \leq \pm 2$ carries a charge $12|y|$ mC/m². Find the total charge on the plate and the electric field intensity at $(0,0,10)$.

$$Q = \int_S \rho_s dS$$

$$Q = \int_{x=-2}^2 \int_{y=-2}^2 12|y| dy dx = \int_{x=-2}^2 dx \int_{y=0}^2 12(2)y dy = 4 \cdot 12(2) \frac{y^2}{2} \Big|_0^2 = 192 \text{ mC}$$

$$\vec{E} = \int \frac{\rho_s}{4\pi\epsilon_0 r^2} dS \hat{a}_r = \int \frac{\rho_s dS}{4\pi\epsilon_0 |\vec{r} - \vec{r}'|^3} (\vec{r} - \vec{r}')$$

$$\vec{r} - \vec{r}' = (0,0,10) - (x',y',0) = (-x',-y',10)$$

Cont... sheet of charge

$$\vec{E} = \int_{-2}^2 \int_{-2}^2 \frac{12|y| dx dy}{4\pi\epsilon_0} \frac{(-x, -y, 10)}{(x^2 + y^2 + 100)^{3/2}}$$

$$= 108 \cdot 10^6 \left[\int_{-2}^2 \int_{-2}^2 \frac{-x dx dy \hat{a}_x}{(x^2 + y^2 + 100)^{3/2}} + \int_{-2}^2 \int_{-2}^2 \frac{-y dy dx \hat{a}_y}{(x^2 + y^2 + 100)^{3/2}} + \int_{-2}^2 \int_{-2}^2 \frac{10|y| dx dy \hat{a}_z}{(x^2 + y^2 + 100)^{3/2}} \right]$$

Due to symmetry only E_z survives:

$$\vec{E} = 108 \cdot 10^6 \cdot 2 \left[\int_{-2}^2 \int_{-2}^2 \frac{10y dx dy \hat{a}_z}{(x^2 + y^2 + 100)^{3/2}} \right]$$

$$= 16.5 \hat{a}_z \text{ MV/m}$$

Chapter Outline

- Coulomb's Law - $F = \frac{kQ_1Q_2}{R^2}$
 - Use when charge distribution is known
- Gauss' s Law - $\Psi = \oint_S \vec{D} \cdot d\vec{S} = Q_{enc}$
 - Use when charge distribution is symmetrical
- Electric Potential (uses scalar, not vectors) $V(r) = -\int_{\infty}^r \vec{E} \cdot d\vec{l}$
 - Use when potential V is known

Electric Flux Density

D is independent of the medium in which the charge is placed.

$$\vec{D} = \epsilon_0 \vec{E} = \int \frac{\rho_v dv}{4\pi R^2} \hat{a}_R \quad [C/m^2]$$

Gauss' s Law

$$\Psi = \oint_S \vec{D} \cdot d\vec{S} = Q_{enc}$$

Therefore:

$$\rho_v = \nabla \cdot \vec{D}$$

This is the 1st of the Maxwell's equations derived here.

Gauss' s Law

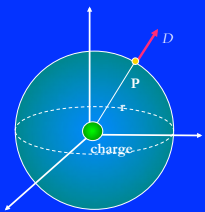
- The total electric flux Ψ through any closed surface is equal to the total charge enclosed by that surface.

$$\vec{D} = \epsilon_0 \vec{E} = \int \frac{\rho_v dv}{4\pi R^2} \hat{a}_R$$

$$\Psi = Q_{enc} = \oint_S \vec{D} \cdot d\vec{S} = \int_V \rho_v dv$$

The key is to choose the Gauss surface to simplify the problem.
Follow the symmetry of the particular case.
Pick surface so that D is \perp

Some examples: Finding D at point P from the charges:



Point Charge is at the origin.

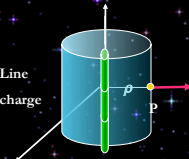
$$Q = \oint_S \vec{D} \cdot d\vec{S}$$

- Choose a spherical dS
- Note where D is perpendicular to this surface.

$$Q = D_r \int_S dS = D_r 4\pi r^2$$

$$D = \frac{Q}{4\pi r^2} \hat{a}_r$$

Some examples: Finding D at point P from the charges:



Infinite Line Charge

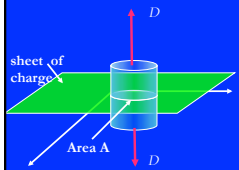
$$\rho_l dl = Q = \oint_S \vec{D} \cdot d\vec{S}$$

- Choose a cylindrical dS
- Note that integral = 0 at top and bottom surfaces of cylinder

$$\vec{D} = \frac{\rho_L}{2\pi\rho} \hat{a}_\rho$$

$$Q = D_\rho \int_S dS = D_\rho 2\pi\rho l$$

Some examples: Find D at point P from the charges:



Infinite Sheet of charge

$$\rho_s \int_S dS = Q = \oint_S \vec{D} \cdot d\vec{S}$$

- Choose a cylindrical box cutting the sheet

$$\rho_s A = Q = D_s \left[\int_{top} dS + \int_{bottom} dS \right]$$

Note that D is parallel to the sides of the box.

$$\rho_s A = D_s [A + A]$$

$$\vec{D} = \frac{\rho_s}{2} \hat{a}_z$$

P.E. A point charge of 30nC is located at the origin, while plane $y=3$ carries charge $10\text{nC}/\text{m}^2$.

4.7 Find D at (0, 4, 3)

$$\vec{D} = \vec{D}_Q + \vec{D}_\rho = \frac{Q}{4\pi r^2} \hat{a}_r + \frac{\rho_s}{2} \hat{a}_n$$

$$\vec{D} = \frac{30 \cdot 10^{-9}}{4\pi(\sqrt{4^2 + 3^2})^2} [(0,4,3) - (0,0,0)] + \frac{10n}{2} \hat{a}_y$$

$$\vec{D} = \frac{30 \cdot 10^{-9}}{4\pi(5)^3} (0,4,3) + 5n\hat{a}_y$$

$$= 5.08\hat{a}_y + 0.057\hat{a}_z \text{ nC/m}^2$$

P.E. 4.8

If $\vec{D} = (2y^2 + z)\hat{x} + 4xy\hat{y} + x\hat{z}$ C/m². **Find :**

- volume charge density at (-1,0,3)
 $\rho_v(-1,0,3) = \nabla \cdot \vec{D} = 4x = -4$ C/m³
- Flux thru the cube defined by $0 \leq x \leq 1, 0 \leq y \leq 1, 0 \leq z \leq 1$

$$\Psi = Q_{enc} = \int_v \rho_v dv = \int_0^1 \int_0^1 \int_0^1 4x dx dy dz$$

- Total charge enclosed by the cube

$$Q = \Psi = 2C$$

Review

Point charge or volume
Charge distribution
$$D = \frac{Q}{4\pi r^2} \hat{a}_r$$

Line charge distribution
$$\vec{D} = \frac{\rho_L}{2\pi\rho} \hat{a}_\rho$$

Sheet charge distribution
$$\vec{D} = \frac{\rho_S}{2} \hat{a}_n$$

We will study Electric charges:

- Coulomb's Law (general cases)
- Gauss' s Law (symmetrical cases)
- **Electric Potential** (uses scalar, not vectors)

Electric Potential, V

- The work done to move a charge Q from A to B is
$$dW = -\vec{F} \cdot d\vec{l} = -Q\vec{E} \cdot d\vec{l}$$
- The (-) means the work is done by an external force.
- The total work = potential energy required in moving Q :
$$W = -Q \int_A^B \vec{E} \cdot d\vec{l}$$
- The energy per unit charge = potential difference between the 2 points:
$$V_{AB} = \frac{W}{Q} = - \int_A^B \vec{E} \cdot d\vec{l} \left[\frac{J}{C} \right] = [V]$$

V is independent of the path taken.

The **Potential** at any point is the potential difference between that point and a chosen reference point at which the potential is zero. (choosing infinity):

$$V(r) = - \int_{\infty}^r \vec{E} \cdot d\vec{l} = - \int_{\infty}^r \frac{Q}{4\pi\epsilon_0 r'^2} \hat{a}_r \cdot dr' \hat{a}_r = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r'} \right]_{\infty}^r = \frac{Q}{4\pi\epsilon_0 r} [V]$$

For many Point charges at r_k :
(apply superposition)
$$V(r) = \frac{1}{4\pi\epsilon_0} \sum_{k=1}^n \frac{Q_k}{|\vec{r} - \vec{r}_k|} [V]$$

For Line Charges:
$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_L \frac{\rho_L(\vec{r}') dl'}{|\vec{r} - \vec{r}'|}$$

For Surface charges:
$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_S \frac{\rho_S(\vec{r}') dS'}{|\vec{r} - \vec{r}'|}$$

For Volume charges:
$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int_v \frac{\rho_v(\vec{r}') dv'}{|\vec{r} - \vec{r}'|}$$

P.E. 4.10 A point charge of $-4\mu C$ is located at $(2,-1,3)$
A point charge of $5\mu C$ is located at $(0,4,-2)$
A point charge of $3\mu C$ is located at the origin

Assume $V(\infty)=0$ and **Find the potential at $(-1, 5, 2)$**

$$V(\vec{r}) = \sum_{k=1}^3 \frac{Q_k}{4\pi\epsilon_0 |\vec{r} - \vec{r}_k|} + C$$

$$|\vec{r} - \vec{r}_1| = (-1,5,2) - (2,-1,3) = \sqrt{46}$$

$$|\vec{r} - \vec{r}_2| = (-1,5,2) - (0,4,-2) = \sqrt{18}$$

$$|\vec{r} - \vec{r}_3| = (-1,5,2) - (0,0,0) = \sqrt{30}$$

$$V(-1,5,2) = \frac{10^{-6}}{1/9 \cdot 10^9} \left[\frac{-4}{\sqrt{46}} + \frac{5}{\sqrt{18}} + \frac{3}{\sqrt{30}} \right] = 10.23 \text{ kV}$$

Example
A line charge of 5nC/m is located on line $x=10, y=20$
Assume $V(0,0,0)=0$ and Find the potential at $A(3, 0, 5)$

$$V(\hat{r}) = -\int \vec{E} \cdot d\vec{l} = -\int \frac{\rho_L}{2\pi\epsilon_0 r} \hat{a}_\rho \cdot d\rho \hat{a}_\rho$$

$$V(\hat{r}) = -\frac{\rho_L}{2\pi\epsilon_0} \ln \rho + C$$

$$V_{origin} - V_A = -\frac{\rho_L}{2\pi\epsilon_0} [\ln \rho_o - \ln \rho_A]$$

$$0 - V_A = -4.8 \quad \quad \quad V_A = +4.8V$$

$\rho_o = |(0,0,0)-(10,20,0)| = 22.36$ and $\rho_A = |(3,0,5)-(10,20,0)| = 21.2$

QUIZ #2: A point charge of 7nC is located at the origin
 $V(0,3,-5)=2V$ and Find C

$$V = \frac{Q}{4\pi\epsilon_0 r} + C$$

P.E. 4.11 A point charge of 5nC is located at the origin
 $V(0,6,-8)=2V$ and Find the potential at $A(-3, 2, 6)$
Find the potential at $B(1,5,7)$, the potential difference V_{AB}

$$V = \frac{Q}{4\pi\epsilon_0 r} + C \quad r = |(0,0,0)-(0,6,-8)| = 10 \quad \therefore C = -2.5$$

$$V_A = \frac{5n}{4\pi\epsilon_0 |(-3,2,6)-(0,0,0)|} + C = 3.93V$$

$$V_B = \frac{5n}{4\pi\epsilon_0 |(1,5,7)-(0,0,0)|} - 2.5 = 2.696V$$


$$V_{AB} = V_B - V_A = -1.233V$$

Relation between E and V
V is independent of the path taken.

$$V_{AB} = -V_{BA}$$

$$V_{AB} + V_{BA} = \oint \vec{E} \cdot d\vec{l} = 0$$

*Esto aplica sólo a campos estáticos.
Significa que no hay trabajo NETO en mover una carga en un paso cerrado donde haya un campo estático E.

$$\oint \vec{E} \cdot d\vec{l} = \int_S (\nabla \times \vec{E}) \cdot d\vec{S} = 0$$


Static E satisfies:

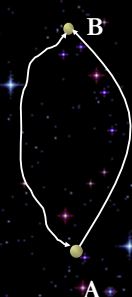
$$\nabla \times \vec{E} = 0$$

Condition for Conservative field = independent of path of integration

$$dV = -\vec{E} \cdot d\vec{l}$$

$$= -E_x dx - E_y dy - E_z dz$$

$$dV = \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz$$

$$\vec{E} = -\nabla V$$


Example Given the potential $V = \frac{10}{r^2} \sin \theta \cos \phi$
Find D at $(2, \frac{\pi}{2}, 0)$.

$$\vec{D} = \epsilon_0 \vec{E} = \epsilon_0 (-\nabla V)$$

In spherical coordinates:

$$\vec{E} = -\left[\frac{\partial V}{\partial r} \hat{a}_r + \frac{1}{r} \frac{\partial V}{\partial \theta} \hat{a}_\theta + \frac{1}{r \sin \theta} \frac{\partial V}{\partial \phi} \hat{a}_\phi \right]$$

$$\vec{E} = -\left[\frac{20}{r^3} \sin \theta \cos \phi \hat{a}_r - \frac{10}{r^3} \cos \theta \cos \phi \hat{a}_\theta + \frac{10}{r^3} \sin \theta \sin \phi \hat{a}_\phi \right]$$

$$\vec{D} = \epsilon_0 \vec{E} \Big|_{(2, \pi/2, 0)} = \epsilon_0 \left(\frac{20}{8} \hat{a}_r + -0 \hat{a}_\theta + 0 \hat{a}_\phi \right)$$

$$\vec{D} = 22.1 \hat{a}_r \text{ C/m}^2$$

P.E. 4.12 Given that $\vec{E} = (3x^2 + y)\vec{a}_x + x\vec{a}_y$ kV/m, find the work done in moving a $-2\mu\text{C}$ charge from $(0,5,0)$ to $(2,-1,0)$ by taking the straight-line path.

a) $(0,5,0) \rightarrow (2,5,0) \rightarrow (2,-1,0)$

$$\frac{-W}{Q} = \int \vec{E} \cdot d\vec{l} = \int [(3x^2 + y)dx + xdy]$$

$$\frac{-W}{Q} = \int_0^2 (3x^2 + 5)dx + \int_5^{-1} xdx$$

$$W = (-Q)(18 - 12) \quad W = 6(-2\mu)12mJ$$

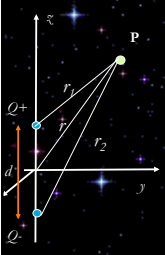
b) $y = 5 - 3x$ $dy = -3dx$

$$\frac{-W}{Q} = \int \vec{E} \cdot d\vec{l} = \int [(3x^2 + 5 - 3x)dx + x(-3dx)] =$$

$$\frac{-W}{Q} = \int_0^2 (3x^2 - 6x + 5)dx \quad \frac{-W}{Q} = 8 - 12 + 10 = 6 \quad W = 12mJ$$

Electric Dipole

- Is formed when 2 point charges of equal but opposite sign are separated by a small distance.



$$V = \frac{Q}{4\pi\epsilon_0} \left[\frac{1}{r_1} - \frac{1}{r_2} \right] = \frac{Q}{4\pi\epsilon_0} \left[\frac{r_2 - r_1}{r_1 r_2} \right]$$

For far away observation points ($r \gg d$):

$$V = \frac{Q}{4\pi\epsilon_0} \frac{d \cos \theta}{r^2}$$

Energy Density in Electrostatic fields

- It can be shown that the total electric work done is:

$$W_E = \frac{1}{2} \int_V \vec{D} \cdot \vec{E} dv = \frac{\epsilon_0}{2} \int_V E^2 dv$$